

The Vibrational Spectra of Strontium Chromate (SrCrO_4) and Lead Chromate (PbCrO_4)

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The Raman spectra of the monoclinic compounds strontium chromate, lead chromate and of the mineral crocoite (PbCrO_4) have been recorded at 300 °K and 77 °K. The spectrum of the orthorhombic modification of lead chromate has also been recorded. The spectra of the monoclinic and of the orthorhombic PbCrO_4 were found to be indistinguishable, whereas the spectra of the isomorphous monoclinic salts SrCrO_4 and PbCrO_4 are remarkably different. Inversion doubling was observed in SrCrO_4 but less distinctly in monocrystalline crocoite.

I. Introduction

SrCrO_4 and PbCrO_4 were investigated as a part of a comprehensive study of the Raman effect in inorganic chromates and dichromates. The particular interest in these substances is in an attempt to shed some light on the influence of the crystal field on the vibrational spectra of anions. It is known that monoclinic SrCrO_4 and PbCrO_4 have almost identical unit cell dimensions. PbCrO_4 exists also in an orthorhombic² modification. No detailed analysis has thus far been undertaken to study the influence of phase transition and isomorphism of chromates by means of the Raman effect.

II. Experimental

1. Apparatus

Raman spectra were recorded on a photoelectrically recording spectrometer described elsewhere^{3,4}. Wave numbers quoted in this paper are accurate to within $\pm 2 \text{ cm}^{-1}$.

For the low temperature work two types of cryostats were used. The first type was a conventional cryostat as used in IR transmission work adapted for 90° scattering experiments by means of a copper cold finger protruding horizontally from the cryostat. A cylindrical Pyrex tube covers the cold finger. This cryostat was used for SrCrO_4 only. The lowest temperature which could be reached with liquid nitrogen as coolant and with the thermocouple embedded in the sample was 110 °K. Since the sample is kept under vacuum and

cooled by conduction only, decomposition of PbCrO_4 by the focussed laser beam ($\lambda = 6328 \text{ Å}$, $P = 180 \text{ mW}$) could not be prevented.

The second type of cryostat simply consists of one Pyrex tube (150 mm \times 10 mm i. d.) surrounded by another (100 mm \times 20 mm o. d.). The space between the tubes, when properly evacuated, serves as a thermal insulation. The sample is put into the centre of the inner tube. Liquid nitrogen is fed into the inner Pyrex tube. A thermocouple pushed in from the other side is used to monitor the temperature. This double-walled Pyrex cryostat has two advantages. Firstly, liquid nitrogen temperature is easily reached and any temperature between 300 °K and 77 °K can be obtained by controlling the N_2 gas flow. Secondly, the sample is cooled by convection at the actual scattering area, so that any heating effect of the absorbed laser radiation is minimised. In the case of PbCrO_4 no decomposition occurs at 77 °K and 160 mW of focussed laser power (power density approximately 10^4 W/cm^2).

2. Sample Preparation

Strontium chromate was prepared by precipitation from aqueous solutions of Merck extra pure potassium chromate and Analar reagent strontium chloride, and dried at 200 °C. The X-ray powder pattern of this material was taken and found to be consistent with the one given by PISTORIUS¹.

Monoclinic lead chromate was prepared by precipitation from aqueous solutions of Merck extra pure potassium chromate and Merck extra pure lead acetate, and dried at 200 °C. The X-ray powder pattern of this compound also was found to be consistent with the one given by PISTORIUS¹.

Lead chromate in its orthorhombic modification was prepared by precipitation from aqueous solutions of

¹ C. W. F. T. PISTORIUS and M. C. PISTORIUS, Z. Krist. **117**, 4 [1962].

² G. COLLOTTI, L. CONTI, and M. ZOCCHI, Acta Cryst. **12**, 416 [1959].

³ W. SCHEUERMANN and G. J. RITTER, Z. Naturforsch. **24 a**, 408 [1969].

⁴ W. SCHEUERMANN and G. J. RITTER, J. Mol. Structure **6**, 240 [1970].



Merck extra pure potassium chromate and Merck extra pure lead nitrate and dried at 25 °C in vacuum. The X-ray powder pattern was found to be consistent with the one given by COLLOTTI, CONTI, and ZOCCHI².

The samples of crocoite of approx. 1 × 1 × 4 mm³ were of Tasmanian origin.

III. Crystal Structure of SrCrO₄ and PbCrO₄

Space groups and unit cells of SrCrO₄ and the two modifications of PbCrO₄ are given in Table 1. The structure of SrCrO₄ has not been determined yet although BRAGG and CLARINGBULL⁵ predicted SrCrO₄ to fall in the barite group which is orthorhombic with space group Pnma – D_{2h}¹⁶, Z = 4. The structure of the mineral crocoite (PbCrO₄) has been established^{6, 7}.

BRODY⁶ could, for crocoite, only determine the positions of the lead and chromium atoms while Quarenì and De Pieri managed to determine the positions of the oxygen atoms and later presented refined data⁷. The structure is as follows. The chromium atom is surrounded by an approximate tetrahedron of oxygen atoms at an average distance of 1.65 Å with bond angles varying from 105.7° to 113.1°. The lead atom is surrounded by seven CrO₄ tetrahedra, with Pb – Cr distances in the range 3.32 Å to 4.15 Å. The lead atom is also surrounded by ten oxygen atoms with Pb – O distances ranging from 2.53 to 3.44 Å; all the other Pb – O distances are greater than 4.03 Å.

The structure of orthorhombic lead chromate PbCrO₄ has also been worked out².

The positions of the atoms were calculated by means of a method based on trial and error. This was done under the assumption of regular CrO₄ tetrahedra. It will be shown in this work that the tetrahedra in SrCrO₄ and both modifications of PbCrO₄ are in fact distorted.

The structure of synthetic monoclinic PbCrO₄ remains to be determined.

IV. Results and Discussion

1. Strontium Chromate

The Raman spectra of strontium chromate at 300 °K and 110 °K were obtained by focussing the laser beam on the powder pellet at an angle of approximately 45° and observing the scattered light perpendicular to the incident beam. Results are listed in Table 2 together with the Raman data of GRIFFITH⁸. The infrared data are taken from AFREMOW and VANDEBERG⁹. The low temperature Raman spectrum is shown in Fig. 1.

A regular CrO₄ tetrahedron possesses *T_d* symmetry. The optically active vibrations are ν_1 : A₁(R), ν_2 : E(R), ν_3 : F₂(R, IR), ν_4 : F₂(R, IR). Neglecting the observed splitting in the first instance, it appears that the spectra can only be explained under the assumption of lifted degeneracies due to a distorted tetrahedron since 7 IR bands and 9 Raman bands have been observed. The total number of 9 Raman bands fits perfectly with the allowed number of vibrations for a five-atomic ion. However, 6 of these are split by 5 to 13 cm⁻¹ in the Raman and 2 bands are split by 10 and 12 cm⁻¹ in the IR. There is a strong indication that the observed splitting is due to inversion doubling of two coupled CrO₄ ions since the splitting of the Raman bands is smaller than that of the corresponding IR bands as demanded by theory¹⁰, viz.,

$$\begin{aligned} \text{R}(857/866 \text{ cm}^{-1}; \Delta\nu &= 9 \text{ cm}^{-1}), \\ \text{IR}(845/855 \text{ cm}^{-1}; \Delta\nu &= 10 \text{ cm}^{-1}), \\ \text{R}(889/893 \text{ cm}^{-1}; \Delta\nu &= 4 \text{ cm}^{-1}), \\ \text{IR}(875/887 \text{ cm}^{-1}; \Delta\nu &= 12 \text{ cm}^{-1}). \end{aligned}$$

Sub- stance	Ref.	System	Space group	Z	a	b	c	β
SrCrO ₄	1	monoclinic	P2 ₁ /n – C _{2h} ⁵	4	7,081	7,388	6,771 Å ± 0,01 Å	103°25' ± 5'
PbCrO ₄	1	monoclinic	P2 ₁ /n – C _{2h} ⁵	4	7,118	7,434	6,794 Å ± 0,004 Å	102°25½' ± 2'
crocoite	7	monoclinic	P2 ₁ /n – C _{2h} ⁵	4	7,12	7,43	6,79 Å	102°25'
PbCrO ₄	2	orthorhombic	Pnma – D _{2h} ¹⁶	4	8,67 ± 0,03	5,59 ± 0,01	7,13 Å ± 0,02 Å	

Table 1. Space groups and unit cells of SrCrO₄ and PbCrO₄.

⁵ L. BRAGG and G. F. CLARINGBULL, *Crystal Structure of Minerals*, Bell, London 1965.

⁶ S. B. BRODY, *J. Chem. Phys.* **10**, 650 [1942].

⁷ S. QUARENÌ and R. DE PIERI, *Acta Cryst.* **19**, 287 [1965].

⁸ W. P. GRIFFITH, *J. Chem. Soc. (A)*, **286** [1970].

⁹ L. C. AFREMOW and J. T. VANDENBERG, *J. Paint Techn.* **38**, (No. 495) 169 [1966].

¹⁰ G. HERZBERG, *Infra-Red and Raman Spectra of Polyatomic Molecules*, D van Nostrand, Inc., New York 1945.

Infra red Ref. ⁹		Raman Ref. ⁸		Powder 300 °K			Powder 110 °K			Assignment
ν	I	ν	I	ν	I	Slit: 2 cm ⁻¹ HW	ν	I	Slit: 1 cm ⁻¹ HW	
[cm ⁻¹]		[cm ⁻¹]		[cm ⁻¹]		[cm ⁻¹]	[cm ⁻¹]		[cm ⁻¹]	T_d
350	vw b	344	(2)	366	116	6.1	337	140	2.4	d, $\nu_2(E)$
		354	(3)	348.7	121	7.6	348	145	2.8	
				354	46	4.2	354	50	4.3	d, $\nu_2(E)$
				367 ^{sh}	40	4.7	367	38	4.7	
410	w sp	380	(6)	373	146	13.2	373	115	8.8	$\nu_4(F_2)$
		407	(4)	401	66	3.8	401	65	2.3	d, $\nu_4(F_2)$
				406	66	3.6	406	72	2.0	
				425	13	3.8	429	15	3.3	d, $\nu_4(F_2)$
431	w sp	437	(2)	432	14	6.6	436	18	3.3	
845	s sp	864	(5)	857	555	5.5	860	750	2.7	d, $\nu_1(A_1)$
855	vw sp sh	870	(10)	866	1030	4.9	867	1250	2.8	
875	vw sp sh			889	600	3.5	892	800	2.1	d, $\nu_3(F_2)$
887	s sp	898	(9)	893	850	4.4	897	950	3.3	
912	s m			915	180	4.9	920	295	2.5	$\nu_3(F_2)$
927	s sp			928	70	4.7	934	100	2.3	$\nu_3(F_2)$

I Intensities are peak values in arbitrary units, cannot be transferred from one spectrum to another and are not corrected for spectral response. d doublet.

Table 2. Vibrational data of SrCrO₄.

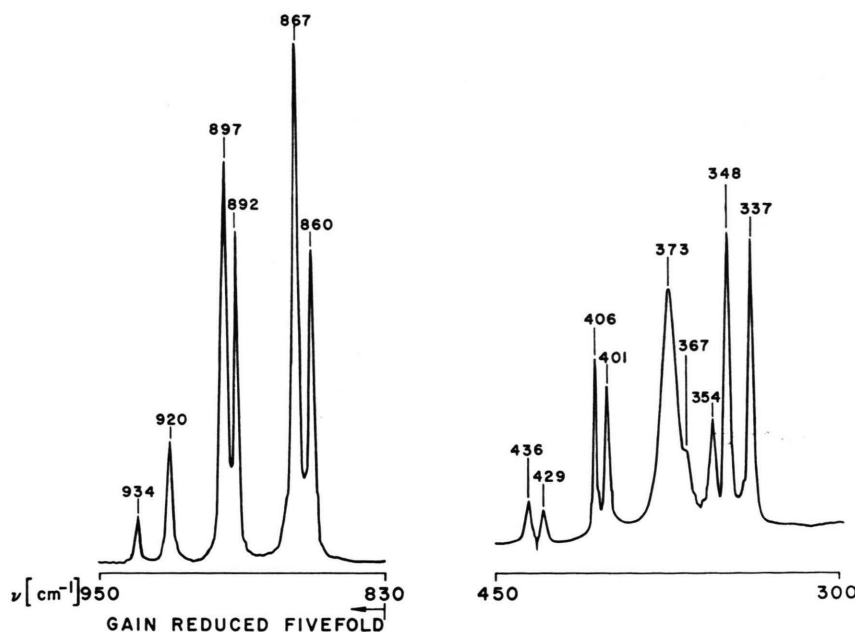
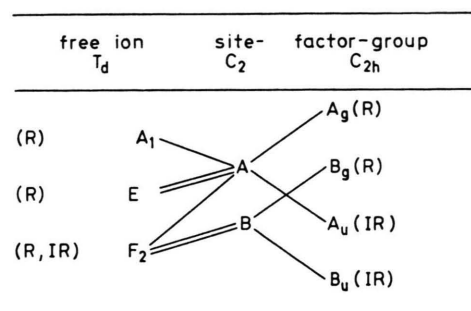


Fig. 1. Low temperature Raman spectrum of SrCrO₄ powder pellet in reflection with 6328 Å/180 mW excitation. $\Theta=110^\circ\text{K}$; $s=50\ \mu=1\ \text{cm}^{-1}$; $T=0.3\ \text{sec}$; $v=5\ \text{\AA}/\text{min}$.

It is reasonable to assume that the strongest band in the Raman spectrum (860/867 cm⁻¹ doublet) is the totally symmetric A₁ vibration. The splitting of this band cannot be explained by means of factor group analysis as can be seen from Fig. 2 where the correlation diagram for the possible site symmetry C₂ is shown.

The other phonons observed in the region between 892 and 934 cm⁻¹ must originate from the triply degenerate $\nu_3(F_2)$ vibration. One component was found split (892/897 cm⁻¹ doublet), the others are single.

The wave number region 337 to 436 cm⁻¹ contains the phonons originating from the $\nu_2(E)$ and

Fig. 2. Correlation diagram for C_{2h} factor group.

$\nu_4(F_2)$ vibrations. Since in tetrahedral ions with a heavy central atom, such as chromium, the doubly degenerate E vibration is usually of lower energy than the triply degenerate F₂ vibration, the two doublets (337/348 and 354/367 cm⁻¹) are assigned to the E vibrations. The remaining phonon at 373 and the doublets at (401/406 and 429/436 cm⁻¹) are assigned to the three F₂ vibrations. The rather large half-width (8.8 cm⁻¹) of the 373 cm⁻¹ band indicates the accidental degeneracy of the two components of one F₂ vibration.

The assignment is listed in Table 2 with reference to a regular tetrahedron.

Another explanation for the observed splittings could be given if two crystallographically and energetically distinct CrO₄²⁻ ions are present in the lattice, giving rise to distinct vibrations. The rule of mutual exclusion would then not be applicable and corresponding bands should occur at the same energies in the infrared and Raman spectra. It can be seen by inspecting Table 2 that this is clearly not the case.

Judging from the vibrational data, it would appear that the structure of SrCrO₄ contains distorted CrO₄²⁻ tetrahedra (symmetry C₂ or C₁) with the Cr atoms in general positions related pairwise by a center of symmetry.

2. Lead Chromate

2.1. Synthetic Lead Chromate

Raman spectra at 300 °K and 77 °K of the monoclinic and at 77 °K of the orthorhombic modification were obtained in the same manner as in the case of SrCrO₄. The results are listed in Table 3 and the low temperature spectrum is shown in Fig. 3.

Although monoclinic strontium and lead chromates belong to the same space group and possess

PbCrO ₄ infra red Ref. ⁹		Crocoite Raman Ref. ⁸		Powder 300 °K			Powder 77 °K		
ν	<i>I</i>	ν	<i>I</i>	ν	Raman <i>I</i>	Slit: 1 cm ⁻¹ HW [cm ⁻¹]	ν	Raman <i>I</i>	Slit: 1 cm ⁻¹ HW [cm ⁻¹]
[cm ⁻¹]		[cm ⁻¹]		[cm ⁻¹]			[cm ⁻¹]		
		180	(1)						
		324	(2)	325	30	6.5	325	48	2.4
		327	(3)						
		337	(5)	337	58	6.5	337	124	2.4
		355	(3)	347	24	5.7	348	40	2.6
		360	(3)	358	150	8.5	358	332	3.1
		377	(4)	377	54	6.5	377	112	2.6
397	m vb, sh	398	(2)	400	22	6.5	401	60	2.2
≈ 790	? sp, sh	795	(3)	n.o.			n.o.		
				n.r.			820	200	5.5
		824	(2)	823 sh			829	200	5.5
833	w m	840	(10)	839	470	13	840	2200	2.5
				n.r.			849 sh	200	5.5
				n.r.			853 sh	100	
858	s m, sh	855	(3)	854 sh			859	320	4.3
≈ 905	w vb								

I Intensities are peak values in arbitrary units, cannot be transferred from one spectrum to another and are not corrected for spectral response.

n. o. Not observed.

n. r. Not resolved.

Table 3. Vibrational data of PbCrO₄ monoclinic.

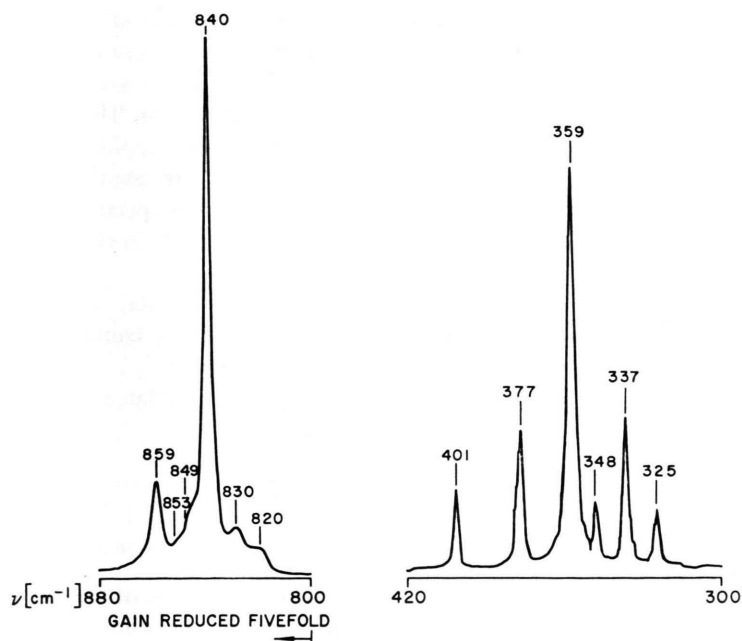


Fig. 3. Low temperature Raman spectrum of PbCrO_4 powder pellet in reflection with $6328 \text{ \AA}/180 \text{ mW}$ excitation. $\Theta = 77^\circ \text{K}$; $s = 50 \mu = 1 \text{ cm}^{-1}$; $T = 0.3 \text{ sec}$; $v = 5 \text{ \AA/min}$.

almost identical unit cells, their vibrational spectra are quite different. On the other hand, the spectra of monoclinic crocoite, monoclinic synthetic lead chromate and orthorhombic synthetic lead chromate are almost indistinguishable. The only significant difference in the Raman spectra of monoclinic and orthorhombic lead chromate lies in the scattering efficiency. The bands of the monoclinic form are approximately eight times as intense as those of the orthorhombic form. This can partly be due to the influence of different particle size in the sample.

The total number of 12 bands in the R and of 5 bands in the IR spectra suggests again that the CrO_4^{2-} ion is a distorted tetrahedron causing all degeneracies to be lifted.

2.2. Crocoite, PbCrO_4

Single crystal Raman spectra of this mineral were taken at 300°K and 77°K . In order to obtain meaningful results¹¹ from a monoclinic crystal with an extraordinary high birefringence ($N_x = 2.31$ Li, $N_y = 2.37$, $N_z = 2.66$; optic plane: 010 ; $Z \wedge c = -5.5^\circ$)¹², it is necessary to choose an arrangement

where neither the incident nor the scattered radiation becomes depolarized. Therefore, the crocoite specimen was cut normal to the c axis and polished to 1 mm thickness. The crystal was then examined in transmission with the incident and scattered radiation parallel to the c axis. This setup was checked by determining the depolarization ratio of the Rayleigh line; it was found to be $\rho_s = 0.01$, with the collecting lens set at 1 : 4 aperture for this and all the single crystal measurements.

The results of these single crystal measurements are listed in Table 4 together with the assignments. The first column gives the wave number of the observed bands, the following four columns give the relative intensities and half width in four meaningful scattering-configurations. The symbol $c(b \perp b)c$ indicates¹³ that the incident and scattered radiation propagate along the c axis, the incident radiation being polarized parallel to the b axis and the scattered radiation being analysed normal to the b axis. It should be noted that the direction $\perp b$ is not coincident with the a axis.

The significance of the low temperature single crystal spectrum can be seen in the excellent resolu-

¹¹ I. R. BEATTIE and T. R. GILSON, Proc. Roy. Soc. London A **307**, 407 [1968].

¹² A. N. WINCHELL and H. WINCHELL, Elements of Optical Mineralogy, Part II, Description of Minerals, 4-th Ed., J. Wiley, New York 1956.

¹³ T. C. DAMEN, S. P. S. PORTO, and B. TELL, Phys. Rev. **142**, 570 [1966].

ν [cm ⁻¹]		$c(b\ b)c$				$c(\perp\ b\ b)c$				$c(\perp\ b\ \perp b)c$				$c(b\ \perp\ b)c$				Assignment	
300 °K	77 °K	300 °K	77 °K		300 °K	77 °K		300 °K	77 °K		300 °K	77 °K		300 °K	77 °K		C _{2h}	T _d	
		<i>I</i>	HW	<i>I</i>	HW	<i>I</i>	HW	<i>I</i>	HW	<i>I</i>	HW	<i>I</i>	HW	<i>I</i>	HW	<i>I</i>	HW		
325	325	4	*	20	3.3	93	6.1	250	2.8	3	*	15	3.3	35	5.2	300	3.0	Bg	E
337	337	133	5.2	280	3.2	6	*	20	3.0	81	5.1	470	3.0	5	*	25	3.3	Ag	E
347	348	52	5.7	140	2.8	58	4.8	180	2.8	17	5.2	40	3.8	21	4.0	170	2.6	Ag + Bg	E
358	358	70	7	250	3.9	6	*	10	*	15	6.1	70	4.3	*	*	5	*	Ag	F ₂
377	377	12	5.7	30	3.8	29	4.5	100	2.8	17	4.3	110	3.0	9	3.9	65	3.0	Ag + Bg	F ₂
400	401	58	4.5	220	2.8	33	4.7	110	2.8	6	4.8	35	4.3	9	4.0	90	2.8	Ag + Bg	F ₂
+	820			160	4.3			45	4.9			360	3.7			25	4.3	Ag	F ₂
823 sh	829	70	+	180	4.9	10	+	15	4.4	46	10	60	3.2	*	*	*	*	Ag	F ₂
839	840	1270	8.5	7000	2.7	70	9	380	2.6	145	11	1000	2.9	29	10	200	2.8	Ag	A ₁
+	847			60	+			150	3.7			30	+			85	4.3	Bg	F ₂
854 sh	859	115	+	540	3.2	46	14	170	4.2	133	13	1400	3.2	19	13	90	4.6	Ag	F ₁
863 sh	867	20	+	30	*	40	12	220	7.5	15	+	30	*	19	12	170	7.1	Bg	F ₂
880 sh	883	*	*	*	*	25	11	20	5.3	*	*	*	*	9	11	20	5.3	Bg	F ₂

I Peak intensities in arbitrary units; same scale for 300 °K and 77 °K spectra; not corrected for spectral response; accuracy $\pm 20\%$.

HW Half width [cm⁻¹] measured with 2 cm⁻¹ spectral slit width; accuracy $\pm 20\%$.

sh Shoulder.

* Too weak to be measured; + too overlapped to be measured.

Table 4. Single crystal Raman data of Crocoite PbCrO₄.

tion obtained, especially in the region from 820 to 883 cm⁻¹. At 300 °K all bands marked sh had to be deduced from their envelopes by graphical methods. The intensities and half width of these bands are, therefore, subject to a larger error.

Considering the Raman tensor of the monoclinic class 2/m – C_{2h} given by LOUDON¹⁴ the assignment of most of the bands is straightforward. Only a few bands show mixed A_g + B_g character.

The connection of the assignment obtained from single crystal measurements with the vibrations of the distorted tetrahedron is done on the basis of the T_d point group. The most intense band in the 820 to 883 cm⁻¹ region is assigned to the ν_1 (A₁) stretching vibration. The only other vibration which can be possibly active in this region is the ν_3 (F₂) bending vibration which is triply degenerate. In order to account for the observed three A_g and three B_g phonons, it must be concluded that the degeneracy of the ν_3 (F₂) vibration is lifted and that each component is split into an A_g and B_g component.

In the region from 325 to 401 cm⁻¹ the doubly degenerate ν_2 (E) bending and the triply degenerate ν_4 (F₂) bending vibration can be expected. Since nine phonons are observed, if the three bands showing mixed character are counted as six separate phonons, the assignment follows the same pattern

as above. All degeneracies are lifted and each component appears split into A_g and B_g. One B_g phonon remains unobserved. The division of the observed phonons into a group originating from E and F₂ vibrations, respectively, can only be made tentatively under the assumption that the E vibration is of lower energy than the F₂ vibration.

3. Inter-ionic Forces

It is interesting to note that the wave number of the totally symmetric stretching vibration ν_1 (A₁) is significantly lower in PbCrO₄ (840 cm⁻¹) than in SrCrO₄ (860/867 cm⁻¹). The weakening of the chromium-oxygen bond can be understood under the assumption of a strengthened lead-oxygen bond which means that the character of the lead-oxygen bond has become partly covalent.

If one allows the insolubility of a chromate to be considered as a measure of bond strength between cation and anion, the conclusion may follow that for the chromate with the highest solubility the wave number of ν_1 (A₁) is also highest and the character of the cation-anion bond has greater ionic character.

In Table 5 wave numbers of vibrations ν_1 (A₁) are listed for a number of chromates, together with the solubilities of these compounds in water¹⁵.

¹⁴ R. LOUDON, Adv. Phys. **13**, 423 [1964].

¹⁵ Handbook of Chemistry and Physics, 40th ed., Chemical Rubber Publishing Co., Cleveland, Ohio 1958.

¹⁶ J. S. STEPHENS and D. W. J. CRUICKSHANK, Acta Cryst. B **26**, 222 [1970].

	CaCrO ₄ · 2 H ₂ O	SrCrO ₄	BaCrO ₄	PbCrO ₄
$\nu_1(A_1)$ at 300 °K	880	857/866	863	839
solubility [g/100 ml] at °C	22.4°	0,12 ¹⁵	0,00034 ¹⁶	0,0000058 ²⁵

Table 5. Wave number of $\nu_1(A_1)$ versus solubility at °C for various chromates.

It can also be argued that the deviation¹⁶ of the Cr—O—Cr bond angle from the ideal value 109.5° observed in crocoite originates from a strengthened Pb—O bond thus causing a decreased bonding power of the narrow tetrahedral d³s orbitals of chromium. This, in turn, would lead to a decrease in the $\nu_1(A_1)$ frequency as observed.

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Partielle Interferenz- und partielle Atomverteilungsfunktionen sowie elektrischer Widerstand von geschmolzenen Magnesium—Zinn-Legierungen *

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From measured intensity functions, published in a previous paper, partial interference functions, partial pair correlation functions, partial atomic distribution functions, and partial coordination numbers were calculated and discussed. According to the theory of Faber and Ziman, the electrical resistivity of molten Mg—Sn alloys was calculated from the partial interference functions and from pseudopotentials reported by Animalu and Heine. The calculated values, being normalized to the resistivity of the molten components, are in good agreement with values measured by a rotating field method.

The partial interference functions are derived by the assumption of independency of concentration. This assumption is proved by discussion of the partial coordination numbers and the results of electrical resistivity determinations.

Für die Berechnungen des elektrischen Widerstandes geschmolzener Legierungen aus Röntgen-Beugungsdaten nach FABER und ZIMAN¹ ist neben der Kenntnis der Pseudopotentiale² die der partiellen Interferenzfunktionen notwendig. Die partiellen Interferenzfunktionen wurden bisher für die Systeme Cu—Sn³, Ag—Mg⁴, sowie Au—Sn und Ag—Sn^{5–7} angegeben. In vorliegender Arbeit werden die partiellen Interferenzfunktionen für das System Mg—Sn aus

den in⁸ publizierten Röntgen-Beugungsexperimenten berechnet. Außerdem werden durch Fouriertransformation der partiellen Interferenzfunktionen die partiellen Atomverteilungsfunktionen und Teilkordinationszahlen berechnet. Mit Hilfe der Interferenzfunktionen erfolgt die Berechnung des elektrischen Widerstandes, der außerdem mit direkt gemessenen Werten⁸ verglichen wird.

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¹ T. E. FABER u. J. M. ZIMAN, Phil. Mag. **11**, 153 [1965].

² A. O. E. ANIMALU u. V. HEINE, Phil. Mag. **12**, 1249 [1965].

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